Petersburg Nuclear Physics

Institute

### **Bonn-Gatchina partial wave analysis:** single energy fit

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#### Main approach: The energy dependent analysis of the data

In many cases an unambiguous partial wave decomposition at fixed energies is impossible. Then the energy and angular parts should be analyzed together:

$$A(s,t) = \sum_{\beta\beta' n} A_{n}^{\beta\beta'}(s) Q_{\mu_{1}...\mu_{n}}^{(\beta)+} F_{\nu_{1}...\nu_{n}}^{\mu_{1}...\mu_{n}} Q_{\nu_{1}...\nu_{n}}^{(\beta')}$$

- 1. Correlations between angular part and energy part are under control.
- 2. Unitarity and analyticity can be introduced from the beginning.
- 3. Three or four body final state + FSI (triangle and box diagrams)
- 4. However, to fix simultaneously energy and angular dependencies of the amplitude a combined fit of many reactions is needed.

# Single energy fit

- 1. First step: decomposition into partial amplitudes (model independent) Limitations:
  - (a) The unique solution demands high precision polarization data.
  - (b) The number of polarization measurements should correspond to a complete experiment.
  - (c) Very difficult to apply to reactions with three or more final states. But:
    - i. Unique decomposition of partial waves.
    - ii. In many cases the resonance contributions are clearly seen.
- 2. Next step: analysis of the partial amplitudes and search for resonances (model dependent)

### Photoproduction amplitude in c.m.s. of the reaction

$$A = \sum_{i} u(k_1) V_{\alpha_1 \dots \alpha_n}^{*(i\pm)\mu} F_{\alpha_1 \dots \alpha_n}^{\beta_1 \dots \beta_n} N_{\beta_1 \dots \beta_n}^{(\pm)} u(q_1) \varepsilon_{\mu} B W_L^{\pm}(s) = \omega^* J_{\mu} \varepsilon_{\mu} \omega' ,$$

$$J_{\mu} = i\mathcal{F}_1\sigma_{\mu} + \mathcal{F}_2(\vec{\sigma}\vec{q})\frac{\varepsilon_{\mu ij}\sigma_i k_j}{|\vec{k}||\vec{q}|} + i\mathcal{F}_3\frac{(\vec{\sigma}\vec{k})}{|\vec{k}||\vec{q}|}q_{\mu} + i\mathcal{F}_4\frac{(\vec{\sigma}\vec{q})}{\vec{q}^2}q_{\mu} .$$

$$\mathcal{F}_{1}(s,z) = \sum_{L=0}^{\infty} [LM_{L}^{+} + E_{L}^{+}]P_{L+1}'(z) + [(L+1)M_{L}^{-} + E_{L}^{-}]P_{L-1}'(z) ,$$
  

$$\mathcal{F}_{2}(s,z) = \sum_{L=1}^{\infty} [(L+1)M_{L}^{+} + LM_{L}^{-}]P_{L}'(z) ,$$
  

$$\mathcal{F}_{3}(s,z) = \sum_{L=1}^{\infty} [E_{L}^{+} - M_{L}^{+}]P_{L+1}''(z) + [E_{L}^{-} + M_{L}^{-}]P_{L-1}''(z) ,$$
  

$$\mathcal{F}_{4}(s,z) = \sum_{L=2}^{\infty} [M_{L}^{+} - E_{L}^{+} - M_{L}^{-} - E_{L}^{-}]P_{L}''(z).$$

#### 1. Reconstruction 4 complex $F_i(s, z)$ from observables

$$Observables \sim \sum_{ij} \mathcal{F}_i \mathcal{F}_j^*$$

at least 8 observables are needed to find  $F_i(s,z)$ , BUT only up to factor  $e^{i\phi(s,z)}$ 

$$E_L^{\pm}(s) = \int_{-1}^{1} \frac{dz}{2} \sum_m \mathcal{F}_m D_m^{(L\pm)} , \qquad (1)$$

However, there is no way to determine multipoles from  $F_i(s,z)e^{i\phi(s,z)}$ 

$$E_L^{\pm}(s) \neq \int_{-1}^{1} \frac{dz}{2} \sum_m \mathcal{F}_m D_m^{(L\pm)} e^{i\phi(s,z)} , \qquad (2)$$

#### **Scattering of spinless particles**

$$A = \sum_{L} A_L(s)(2L+1)P_L(z)BW_L(s)$$

$$A_{L}(s) = \int_{-1}^{1} \frac{dz}{2} A P_{L}(z)$$

In experiment only differential cross section can be measured, which is proportional to  $|A|^2$ :

$$A_L(s) \neq \int_{-1}^{1} \frac{dz}{2} |A| P_L(z)$$

BUT  $|A|^2 = |A_0|^2 + |A_1|z^2 + \ldots + 2Re(A_0A_1^*)z + \ldots$ 

#### Single energy analysis is a multipole decomposition

Then the following questions should be answered:

- 1. How many multipoles one needs to describe the full set of data (all polarization observables) at certain energy?
- 2. What additional ambiguities appear in multipole decomposition?
- 3. What are effects from statistics and acceptance (e.g. not full angular coverage).

#### **Limitations:**

$$Observables(s, z) \sim \sum \mathcal{M}ultipole_i(s)\mathcal{M}ultipole_j^*(s)D(z)$$

Multipoles can be found up to the arbitrary phase  $\phi(s)$ . In our analysis only absolute value for  $M_1^+$  is fitted.

## Single energy fit program.

- Investigation of the  $\gamma p \to \pi^0 p$  reaction: small contribution from t-channel exchanges, waves with low orbital momentum are dominant at low energies.
  - 1. Fit of pseudo data: observables predicted by the Bonn-Gatchina PWA solution taken with small errors and full angle coverage.
  - 2. Investigation of effects from statistic and acceptance. E.g., realistic errors, not full angular coverage.
- Investigation of the  $\gamma p \to \pi^+ n$  reaction: large contribution from t-channel exchanges, waves with high orbital momentum are important already at low energies.

Same investigation as for  $\gamma p \to \pi^0 p$ 



Description of the  $\gamma p 
ightarrow \pi^0 p$  pseudo data - multipoles up to L=1 and L=2



Description of the  $\gamma p 
ightarrow \pi^0 p$  pseudo data - multipoles up to L=1 and L=2



Description of the  $\gamma p 
ightarrow \pi^0 p$  pseudo data - multipoles up to L=3 and L=4



Description of the  $\gamma p 
ightarrow \pi^0 p$  pseudo data - multipoles up to L=3 and L=4



**Real** and **imaginary** parts in compare with BoGa PWA multipoles



#### **Real** and **imaginary** parts compared with BoGa PWA multipoles







**Real** and **imaginary** parts compared with BoGa PWA multipoles







**Real and imaginary parts compared with BoGa PWA multipoles** 



**Real** and **imaginary** parts compared with BoGa PWA multipoles

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Fit: real and imaginary parts compared with BoGa PWA multipoles



**Real** and **imaginary** parts compared with BoGa PWA multipoles

### Conclusion

- 1. To fit pseudo  $\gamma p \rightarrow \pi^0 p$  data up to  $W \sim 1300$  MeV one needs up to L = 2 multipoles At higher energy up to W = 1600 - 1700 MeV L = 3 multipoles are needed; at W = 1900 MeV L = 4 multipoles are needed
- Reconstructed multipoles agree well with the input multipoles from Bonn-Gatchina PWA. Higher multipoles which are omitted from the single energy fit influence only highest L multipoles taken into account.
- 3. The higher multipoles can be approximated with second order polynomials.

#### **Next steps**

- 1. To investigate effects from acceptance (for  $\gamma p \rightarrow \pi^0 p$ ).
- 2. To start investigation of the  $\gamma p \rightarrow \pi^+ n$  reaction.



Description of the  $\gamma p 
ightarrow \pi^+ n$  pseudo data - multipoles up to L=3













**Real and imaginary parts compared with BoGa PWA multipoles**